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DELAWARE UNIV NEWARK DEPT OF CIVIL ENGINEERING
COMPUTER MODEL FOR ENERGY SPECTRUM TRANSFORMATION OVER IRREGULA--ETC(U)
JUL 76 H WANG, J C SHIAU
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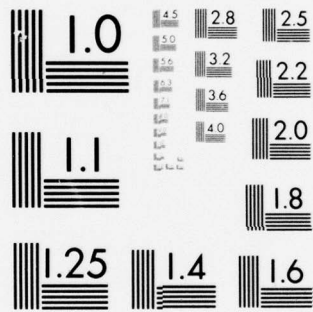
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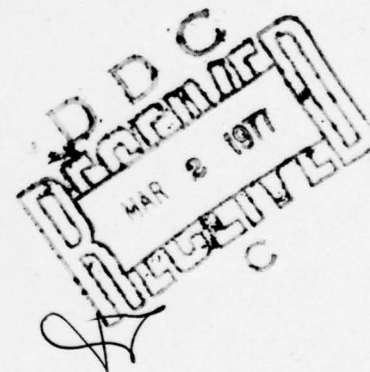
**Computer Model for
Energy Spectrum Transformation
Over Irregular Bottom Topography**

by

Hsiang Wang

and

J. C. Shiau



Technical Report No. 2

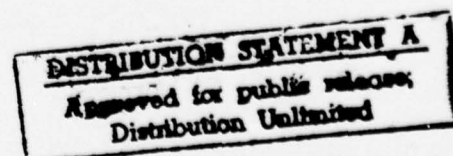
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Department of Civil Engineering
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LIST OF SYMBOLS

A_o	amplitude of shoreline undulation
e	wave phase velocity = L/T
C_g	wave group velocity
E	wave energy
f_i	index column for rhythmic topography
F	spectral density function
g	gravitational acceleration
h	local water depth
h_b	water depth at wave breaking point
h_o	elevation at the origin of coordinate system
H	wave height
i	grid index x-direction
j	grid index y-direction
k	magnitude of the wave number = $2\pi/L$
k_o	deepwater wave number
$K(\theta)$	directional spreading function for wave energy
L	wave length
m	beach slope
r	magnitude of position vector
\vec{r}	position vector
S	extent in offshore direction where bottom undulation begins, beyond which the contour is parallel
$S(\omega)$	one-dimensional spectral density function
t	time

List of Symbols (Continued)

T	wave period
u	wind speed
x	coordinate perpendicular to shoreline
y	coordinate parallel to shoreline
Δx	grid size in x-direction
Δy	grid size in y-direction
α	polar angle of the spectral vector
γ	specific weight of fluid
θ	wave direction
ξ	dummy variable
ϕ	the frequency directional wave spectral density function
ϕ_0	deepwater frequency, directional wave spectral density function
ω	wave angular frequency = $2\pi/T$

ABSTRACT

Based on the assumption that the wave energy associated with a narrow frequency band stayed within the band upon refraction, a numerical model based on finite difference method has been developed. This model, utilizing electronic computer, proves to be quite convenient to determine shallow water wave spectra at designated spatial locations of irregular bottom contours provided the deepwater wave spectrum is given. The model has been verified against analytical solution for case of two-dimensional parallel bottom contours and compared in good agreement with field measurement at location well beyond the surf zone.

CHAPTER I INTRODUCTION

The needs for more realistic representation and accurate prediction of shallow water environment are increasing for both scientific and engineering purposes, in particular, due to the advent of anticipated continental shelf development. It is also an accepted opinion that wave spectrum which utilizes the ensemble of random surface oscillations provides a more realistic, and perhaps more accurate, representation of ocean waves than monochromatic wave train of regular oscillation. In the past, considerable efforts have been devoted in the development of deep water wave spectrum to the extent that one has reasonable confidence on its applicability to describe wind generated ocean waves. The goal of this work is to develop a practical method of transforming this deepwater water wave spectra into shallow water of irregular bottom topography.

A number of investigators (Longuet-Higgins, 1956; Collins, 1970; etc.) have studied the theoretical formulation of the transformation of wind wave spectrum in shoaling water by considering the conservation of energy between adjacent wave rays much the same way of the wave refraction in monochromatic wave train. Based on Longuet-Higgins' formulation, Karlsson (1969) developed a method to compute spectrum transformation over parallel bottom contours. Collins (1972) extended the work to the inclusion of bottom friction and have mentioned the application to the irregular bottom topography. However, his numerical scheme traces wave energy along wave rays which makes the computational procedure quite

impractical in determining shallow water wave spectrum at designated locations over an area of irregular offshore bottom configuration. More recently, Krasitkiy (1974) has derived explicit analytical solutions for the spectrum transformation over two-dimensional parallel bottom contours.

The present study presents a numerical method which proves to be effective to transform deepwater wave spectrum into shallow water regions. It is based on the assumption that the wave energy associated with a narrow frequency band stays within this band upon refraction when waves propagate from deep into shallow waters. A recently developed wave refraction program (Noda, et al., 1974) is modified to compute the spatial modification of the incident wave energy within the designated frequency. The shallow water wave spectrum at any specified location, is then simply the contributions resulting from energy associated with different frequency bands of the spectrum.

Examples of the spectrum distribution over parallel bottom contours, irregular bottom bathymetry and rhythmic topography were illustrated. Comparisons of the present results with analytic solution obtained by Krasitskiy and the field observations were made.

CHAPTER II TRANSFORMATION OF ENERGY SPECTRUM

Let $F(k, r, t)$ be the spectral density function defined as energy per unit area per unit wave number space, the total energy is equal to the summation of the spectral density over the whole wave number space, i.e.,

$$E = \int_0^{\infty} F(k, \vec{r}, t) dk \quad (1)$$

where k = magnitude of the wave number = $2\pi/L$

L = wave length

\vec{r} = position vector

t = time

For the case of steady state, the wave energy contained within a small increment dk can be expressed as

$$\Delta E = \int_k^{k+dk} F(\xi, \vec{r}) d\xi$$

Through Taylor's expansion and neglecting higher orders, we have

$$\Delta E = F(k, \vec{r}) dk \quad (2)$$

It is known that for linear wave systems, the following dispersion relationship exists

$$\omega^2 = gk \tanh kh \quad (3)$$

where ω = wave angular frequency = $2\pi/T$

T = wave period

h = local water depth

g = gravitational acceleration

Differentiating the above equation results in

$$d\omega = C_g dk \quad (4)$$

where C_g = wave group velocity = nc

$$n = (1 + 2kh/\sinh 2kh)/2$$

$$c = \text{wave phase velocity} = L/T$$

Utilizing the above relation, the wave number spectral density function $F(k, \vec{r})$ can be transformed into polar frequency direction space as

$$F(k, r) = \frac{C_g}{k} \phi(\omega, \alpha, r) \quad (5)$$

where ϕ = the frequency directional wave spectral density function

r = magnitude of space vector \vec{r}

α = polar angle of the spectral vector

Combining Equations (2) and (5), one obtains

$$\phi(\omega, \alpha, r) = \frac{k \Delta E}{d\omega} \quad (6)$$

Assuming that the wave energy associated with a narrow frequency band would stay within this band upon wave refraction, a relationship between shallow water and deepwater wave energy spectra can be obtained by

$$\phi = \frac{k}{k_o} \frac{\Delta E}{\Delta E_o} \phi_o \quad (7)$$

where the subscript o denotes deepwater conditions.

For small amplitude waves, it is known that (Lamb, 1932)

$$E = \gamma \cdot \frac{H^2}{8}$$

where γ = specific weight of the fluid

H = wave height

It is plausible to assume here that

$$\Delta E = \gamma \cdot \frac{(\Delta H)^2}{8} \quad (8)$$

where ΔH can be interpreted as the contribution to the total wave height due to the associated increment of wave energy. Based on this relationship, Equation (7) can be rewritten as

$$\phi = \frac{k}{k_o} \frac{(\Delta H)^2}{(\Delta H_o)^2} \phi_o \quad (9)$$

This equation which clearly reveals the effects of shoaling and dispersion is the key equation for wave energy transformation computations. It is evident from this equation that the task of computing ϕ requires the computation of wave height and wave number transformations at designated locations.

CHAPTER III TRANSFORMATION OF WAVE CLIMATE

To compute wave height and wave number transformation at designated locations in shallow water, a numerical program on wave refraction and shoaling recently developed by Noda et.al. (1974) is adopted and modified. This method like most of the existing techniques in computing wave refraction, neglected the nonlinear effects, bottom friction, turbulent dissipation and bottom percolation. The procedure involves solving a pair of simultaneous equations for the wave direction and the wave height. The first equation is derived from the fact that the wave number vector field, k , is irrotational (Phillips, 1966). In Cartesian coordinates as shown in Figure 1, the following equation results

$$\cos\theta \frac{\partial\theta}{\partial x} + \sin\theta \frac{\partial\theta}{\partial y} = \frac{1}{k} (\cos\theta \frac{\partial k}{\partial y} - \sin\theta \frac{\partial k}{\partial x}) \quad (10)$$

where θ = angle of wave incidence to the beach normal

x = Cartesian coordinate perpendicular to the shoreline

y = Cartesian coordinate parallel to the shoreline

The partial derivatives of k were determined from the wave dispersion relationship expressed in Equation (3). The second equation is obtained from the condition of the steady state conservation of energy (Phillips, 1966).

$$\begin{aligned} (C_g \cos\theta) \frac{2}{H} \frac{\partial H}{\partial x} + (C_g \sin\theta) \frac{2}{H} \frac{\partial H}{\partial y} + \cos\theta \frac{\partial C_g}{\partial x} - C_g \sin\theta \frac{\partial\theta}{\partial x} \\ + \sin\theta \frac{\partial C_g}{\partial y} + C_g \cos\theta \frac{\partial\theta}{\partial y} = 0 \end{aligned} \quad (11)$$

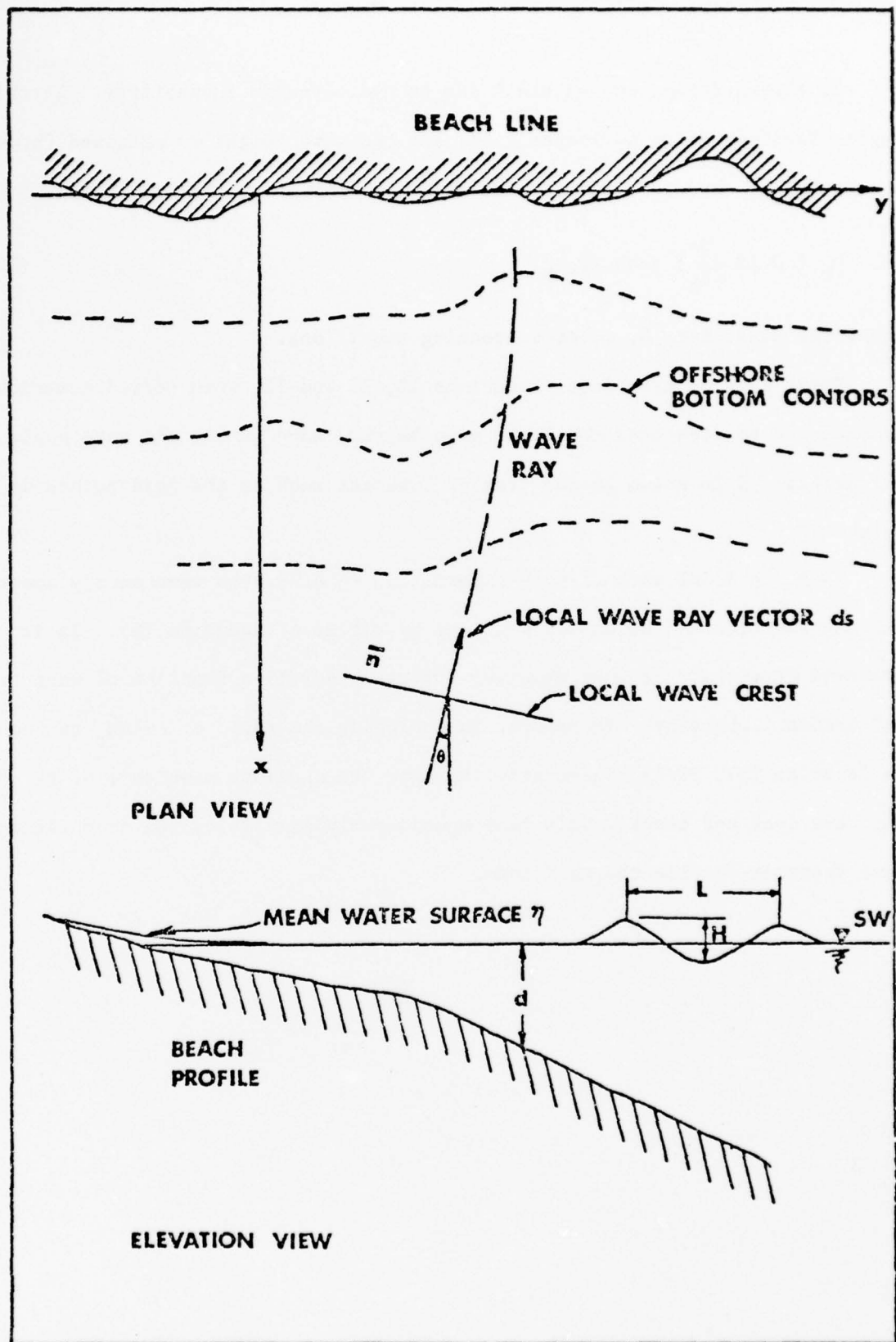


Figure 1 Definition Sketch of Nearshore Environment

In the surf zone, waves break due to hydrodynamic instability. In this region Equation 11 is no longer valid and the wave height is obtained through the use of an empirical equation developed by LeMehaute and Koh (1967).

$$H_b = 0.12 \left(\frac{2\pi}{k_b} \right) \tanh(k_b h_b) \quad (12)$$

where the subscript, b, denotes breaking conditions.

These sets of equations, Equations 10, 11 and 12, when solved numerically, enables one to determine the local wave height, wave number and wave angle at any designated location in the area of interest such as the grid points in Figure 2.

With the local wave climate information in hand, the wave energy spectral density function can be easily obtained by virtue of Equation (9). It is observed here that the wave shoaling effect is purely a function of wave frequency and bottom topography. Therefore, in computing the ratio of $\Delta H / \Delta H_0$ as required in Equation (9), it is nonessential to know the absolute magnitude of ΔH_0 provided the wave does not break. This fact considerably simplifies the computation of wave spectrum outside the surf zone.

CHAPTER IV NUMERICAL PROCEDURES

Computational Scheme. The computation of the transformation of wave climate follows closely the method developed by Noda (1974). Equations of (10) and (11) are written in finite difference form using a forward difference in x, and a backward difference in y for the local grid system shown in Figure 3. Each grid was assigned a mean depth h_{ij} and an area $\Delta x \Delta y$. All major quantities--H, θ , k--are computed at the grid center. Values of $\sin \theta_{i,j}$ and $\cos \theta_{i,j}$ are determined using an average of these quantities from four surrounding grid blocks approximated to second order in a Taylor series expansion, i.e.,

$$\begin{aligned} \sin \theta_{i,j} = & \frac{1}{4} (\sin \theta_{i+1,j} + \sin \theta_{i-1,j} + \sin \theta_{i,j+1} + \sin \theta_{i,j-1}) \\ & + \frac{1}{8} [(\theta_{i+1,j} - \theta_{i-1,j})(\cos \theta_{i-1,j} - \cos \theta_{i+1,j}) + (\theta_{i,j+1} - \theta_{i,j-1}) \\ & (\cos \theta_{i,j+1} - \cos \theta_{i,j-1})] \end{aligned} \quad (13)$$

etc.

The numerical iterations are carried to an acceptable error, in this case,

$$\frac{|X_{\text{new}} - X_{\text{old}}|}{|X_{\text{new}}|} \leq 0.1\% \quad (14)$$

where X represents all the variables under consideration.

The computation of spectrum transformation does not involve numerical integration and is straight forward from Equation (9).

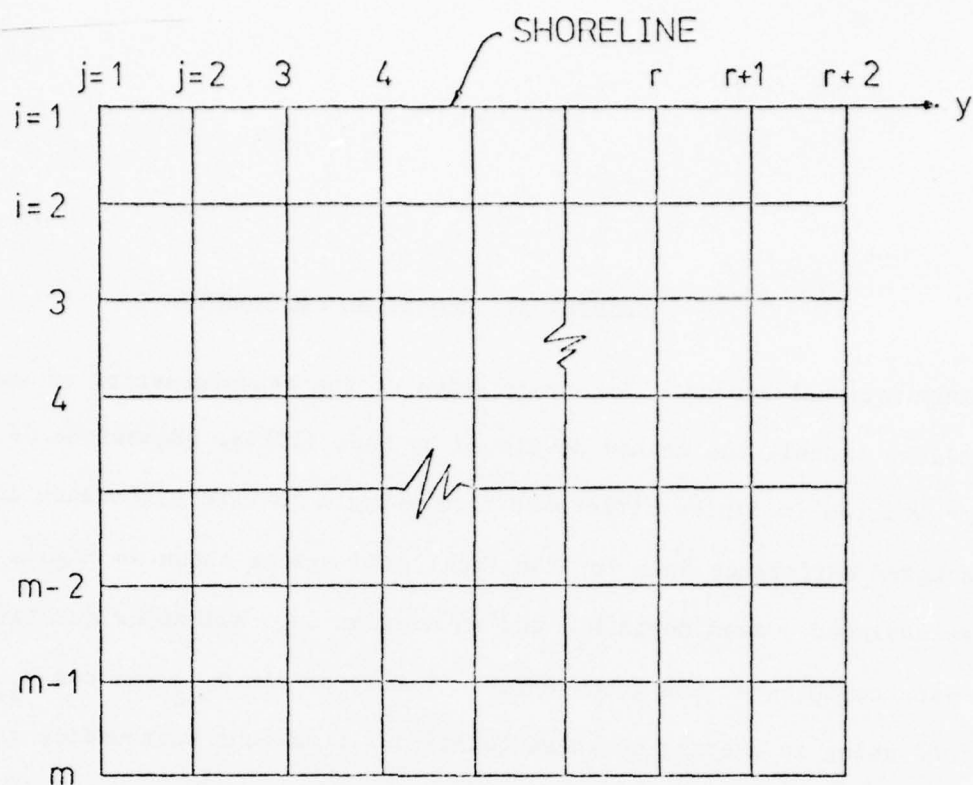


Figure 2 Grid System Used in Computation

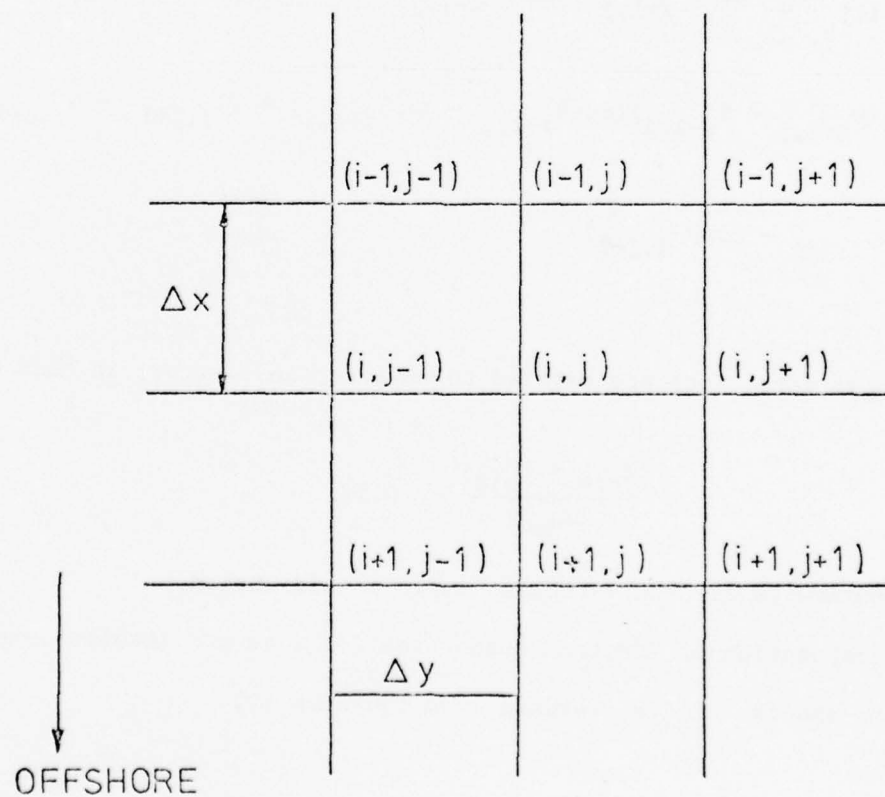


Figure 3 Local Grid Description

Boundary Conditions. The condition at the shoreward boundary is simply that the wave angles are normal to the shoreline at the boundary. Two different options are built in the program for the side boundaries.

A. For rhythmic topography--In this case the wave condition repeats itself at space intervals equal to the integer multiplications of the periodic length. The boundary condition can be self-generated by adding additional grid column such that

$$f_{i, N+1} = f_{i,2} \quad (15)$$

where $N+1$ = periodic length.

A. For slowly varying irregular topography--The conditions at the upwave side are computed by extending the bottom profiles two dimensionally such that the wave refraction follow the Snell law up to the boundary. The downward side boundary condition needs not be prescribed for this case.

Input Wave Condition. Strictly speaking, directional spectrum of the form $\phi_B(\omega, \theta)$ should be prescribed at the seaward boundary. Since our present knowledge on directional spectrum is quite inadequate, the program, therefore, is written to accept the deepwater wave spectrum of the following form as input:

$$\phi_0(\omega, \theta) = S(\omega) K(\theta) \quad (16)$$

where $K(\theta)$ is an artificial weighting function to estimate the effect of directional spreading. For the sample computations to be illustrated $S(\omega)$ is taken to be the Pierson-Moskowitz's deep water wave spectrum:

$$S(\omega) = \alpha \frac{g^2}{\omega^5} \exp \left[-\beta \left(\frac{g}{\omega U} \right)^4 \right] \quad (17)$$

where U is the wind speed, $\alpha = 8.1 \times 10^{-3}$ and $\beta = 0.74$.

The function of $K(\theta)$ assumes the form of

$$K(\theta) = \begin{cases} (8/3\pi) \cos^4 \theta & |\theta| < \pi/2 \\ 0 & |\theta| \geq \pi/2 \end{cases} \quad (18)$$

to facilitate comparisons with analytical results of Krasitskiy (1974).

As the seaward boundary is often set at finite water depths to avoid excessive computational time the deep water wave spectrum is transformed to the finite depths at the boundary using analytical solution of Krasitskiy (1974) with the assumption that the offshore topography can be approximated by two-dimensional parallel contours.

CHAPTER V VERIFICATION AND APPLICATIONS

Sample computations of three different types of bottom hydrographs are illustrated here. The first example is for a two-dimensional parallel contour case such that the results can be compared with analytical solution of Krasitskiy (1974). The calculations are carried out for a wind speed of 20.5m/sec (the same value chosen by Krasitskiy) and a uniform slope $m = 0.005$. Figure 4 compares the numerical results of spectral density function with the analytical ones for water depths of ∞ , 25M and 15M. The present numerical scheme yields satisfactory results except at high frequency end where small deviations from the analytical values are noted. This discrepancy is considered inevitable due to the finite grid size. Reducing grid size will improve the numerical errors at the expense of increasing computational time.

The second case is a comparison between computed spectrum with the field measurements. The spectrum transformation along the Delaware coast, facing the Atlantic Ocean, located approximately two miles north of Indian River Inlet (Figure 5) was examined. Field measurements were obtained by capacitance-type wave gage mounted on a wave tower in water of approximately 4.4m deep. The location of the tower and the bottom bathymetry is shown in Figure 6. Data was recorded on tape. Wave energy spectrum was obtained through fast-Fourier analysis. Detailed descriptions of the field measuring program is documented by

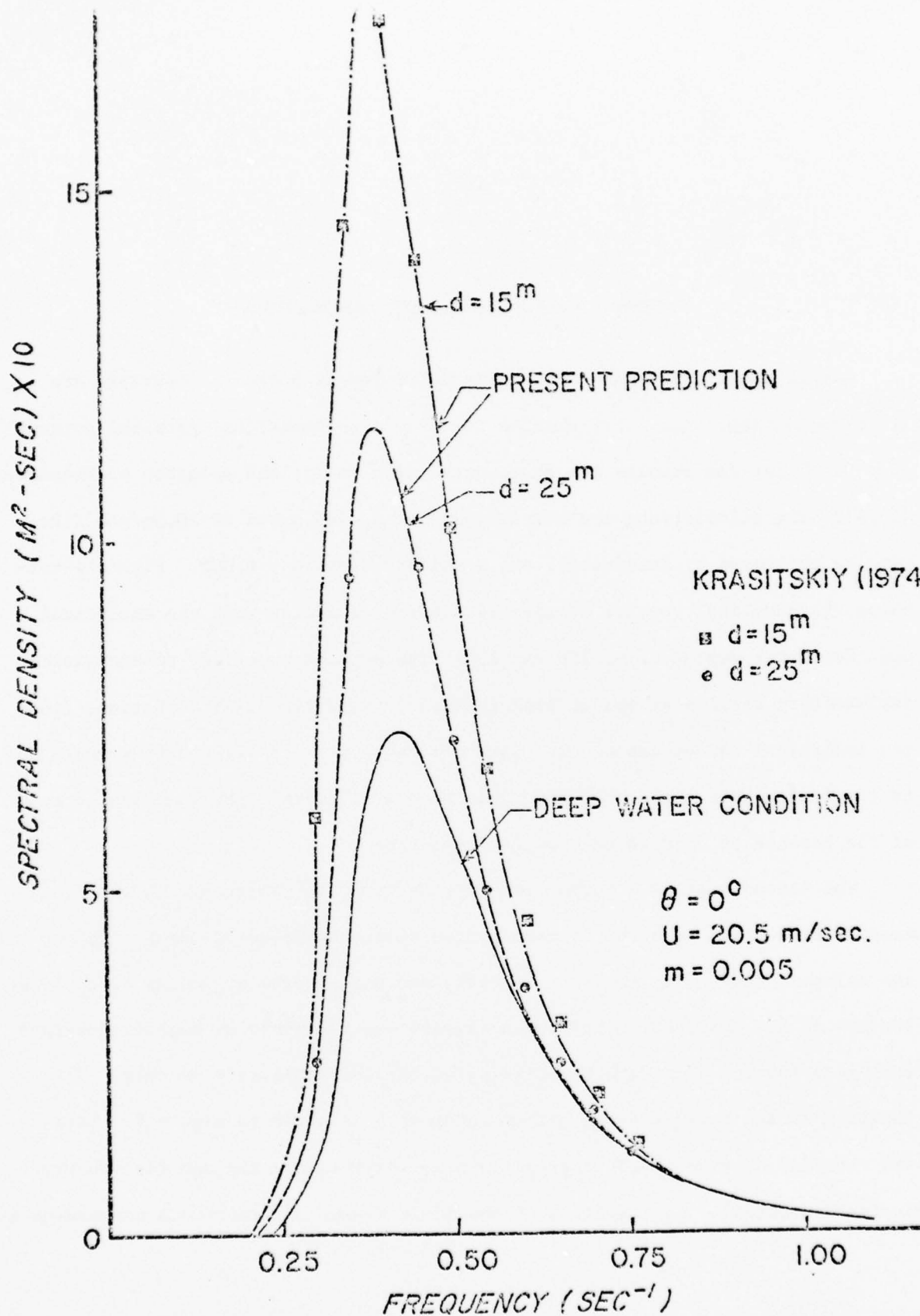


Figure 4 Comparisons of Numerical Results and Analytical Results of Spectrum Transformation Over Parallel Bottom Contour

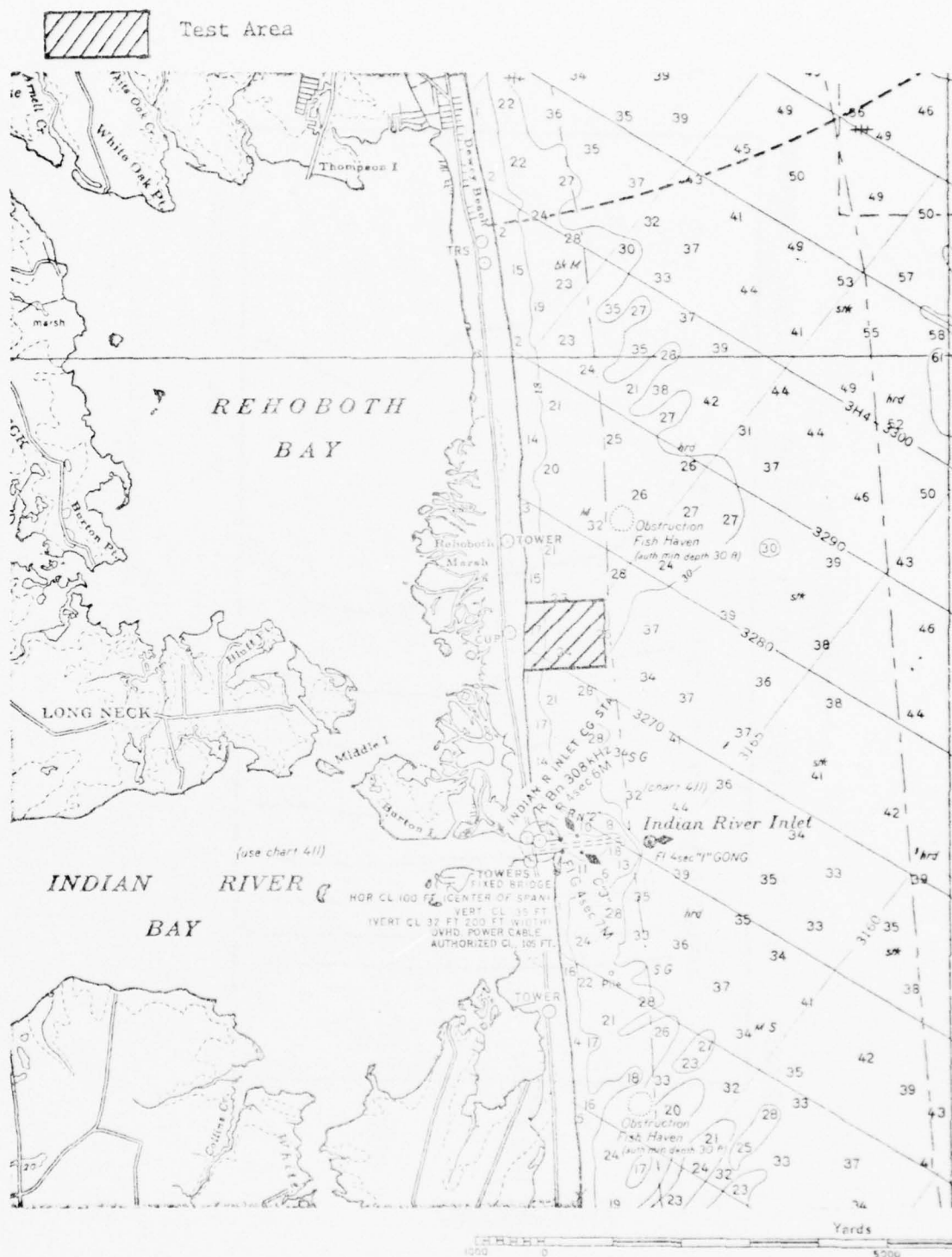


Figure 5 Field Measurement Site, Indian River Inlet, Delaware
(From Map C&GS 1219)

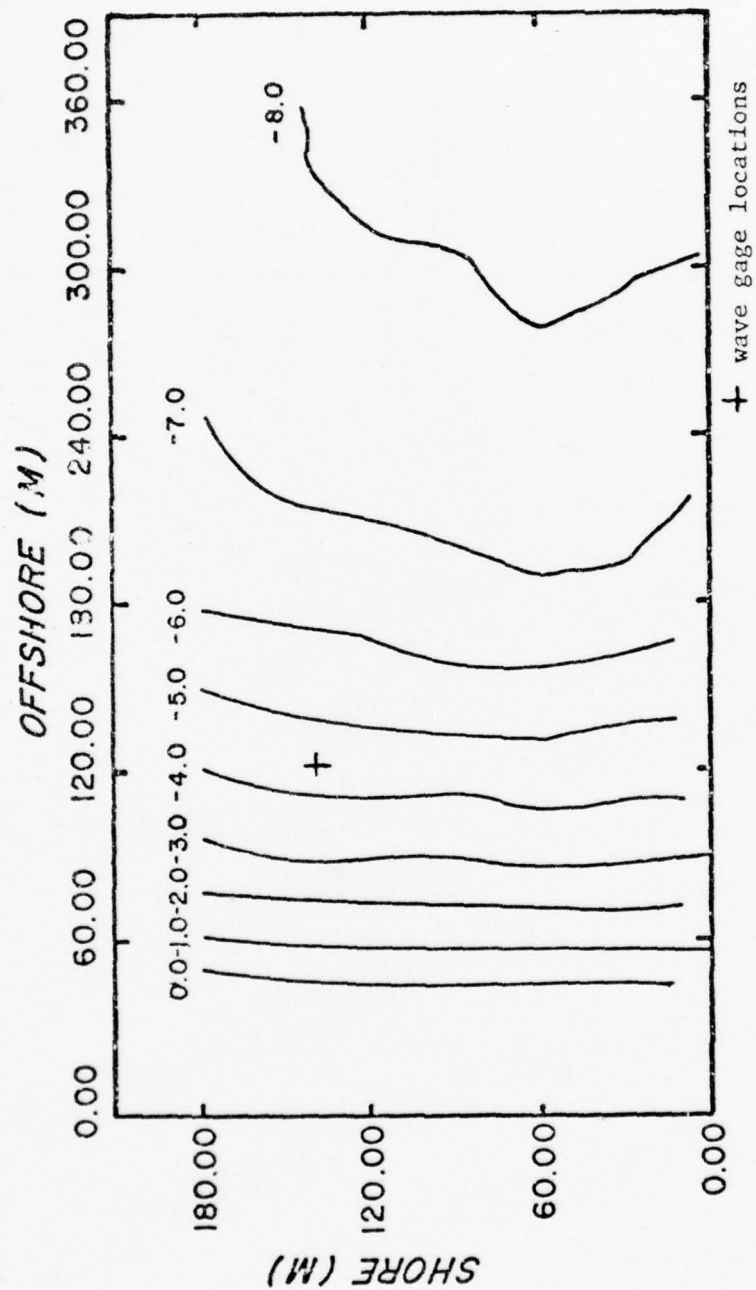


Figure 6 Bottom Profiles and Wave Gage Location

Borchardt (1974) and Hsu (1975). Samples of comparisons are shown in Figure 7. Since there is no simultaneous measurement of deep water wave conditions, the deep water wave spectrum of Pierson-Moskowitz type is again adopted as input condition in the numerical computation. A north-easterly wind of 8.3 m/sec is assumed to prevail. Since the recorded wind speed only indicated a range of 8 - 10 m/sec, the selection of 8.3 m/sec in the computation is merely an estimation. Varying this input wind speed will affect basically the magnitude of the spectral density but hardly the shape of the function.

Finally, the spectrum transformation over a rhythmic nearshore hydrograph such as the type described by Sonu (1972) is demonstrated. Assuming that the bathymetry undulates alongshore with diminishing amplitude toward offshore, as shown in Figure 8, such that the water depth, h , in the x - y plane, can be expressed as the superposition of a uniformly sloped beach and a sinusoidal variation:

$$h = -m(x-\Delta x) + A_o \left(1 - \frac{x}{S}\right) \sin \frac{\pi}{2} \left(1 - \frac{y}{\Delta y}\right) + h_o \quad (20)$$

for $x \leq S$ and

$$h = -m(x-\Delta x) + h_o$$

for $x > S$

where m = beach slope

A_o = amplitude of shoreline undulation

h_o = elevation at origin of coordinate system

S = extent in the offshore direction where bottom undulation begins,
beyond which the contour is parallel.

$\Delta x, \Delta y$ = grid sizes along the x and y directions respectively.

The values used in the numerical examples are: $m = 0.02$, $A_o = 0.13M$, $h_o = 1M$, $S = 360M$, and $\Delta x = \Delta y = 30M$.

The computed spectral density functions at nodal and anti-nodal points are shown in Figure 9. The results clearly indicate, as expected, that energy flux converged at the convex section of the shoreline and diverges along the concave section.

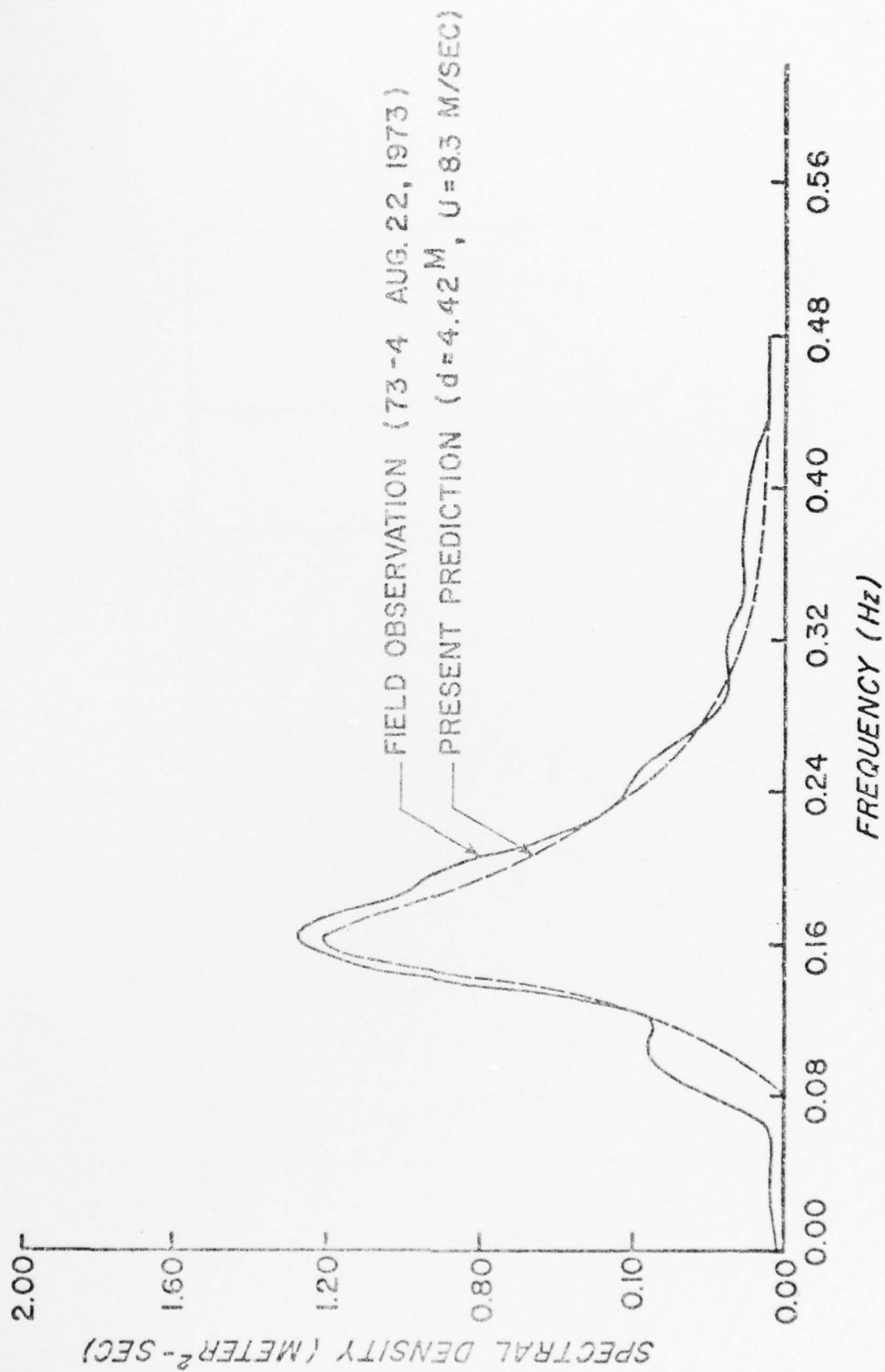


Figure 7 Comparisons of Numerical and Field Results of Spectrum Transformation

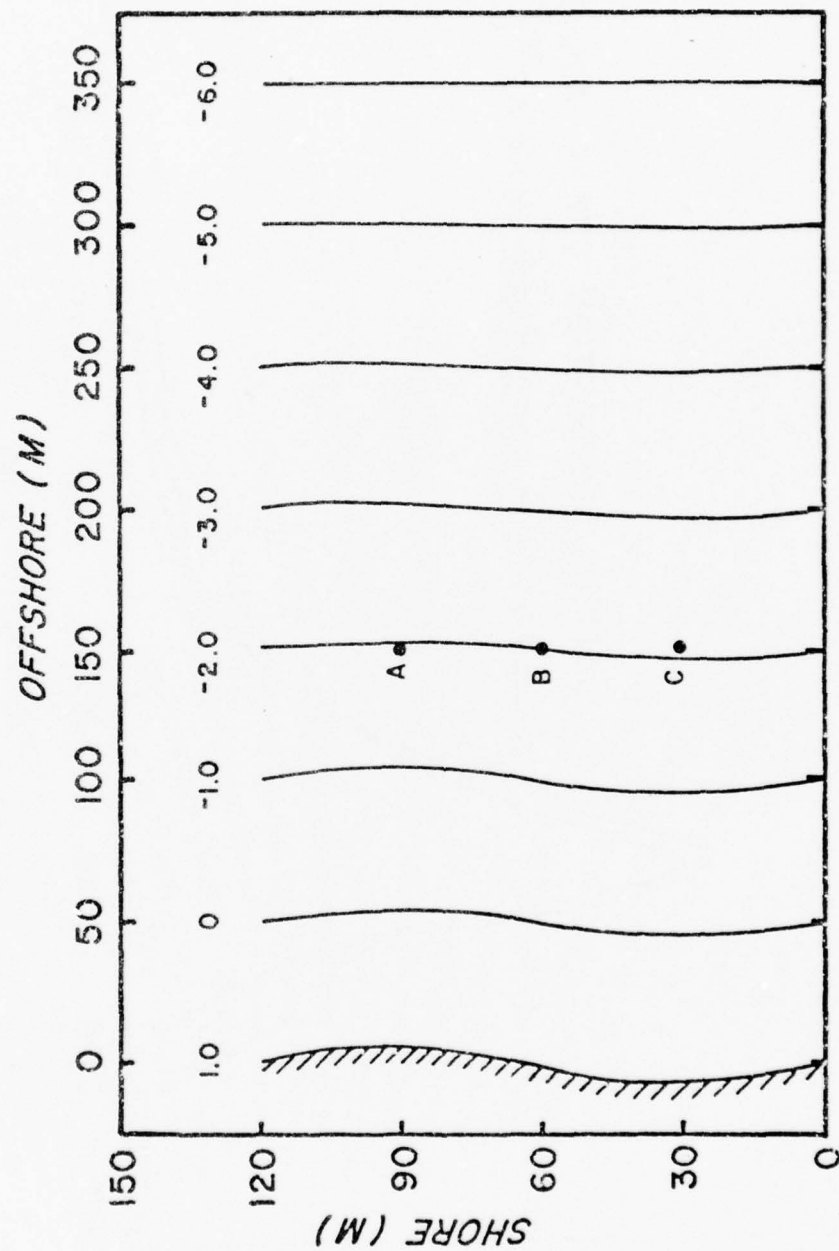


Figure 8 Hypothetical Rhythmic Topography

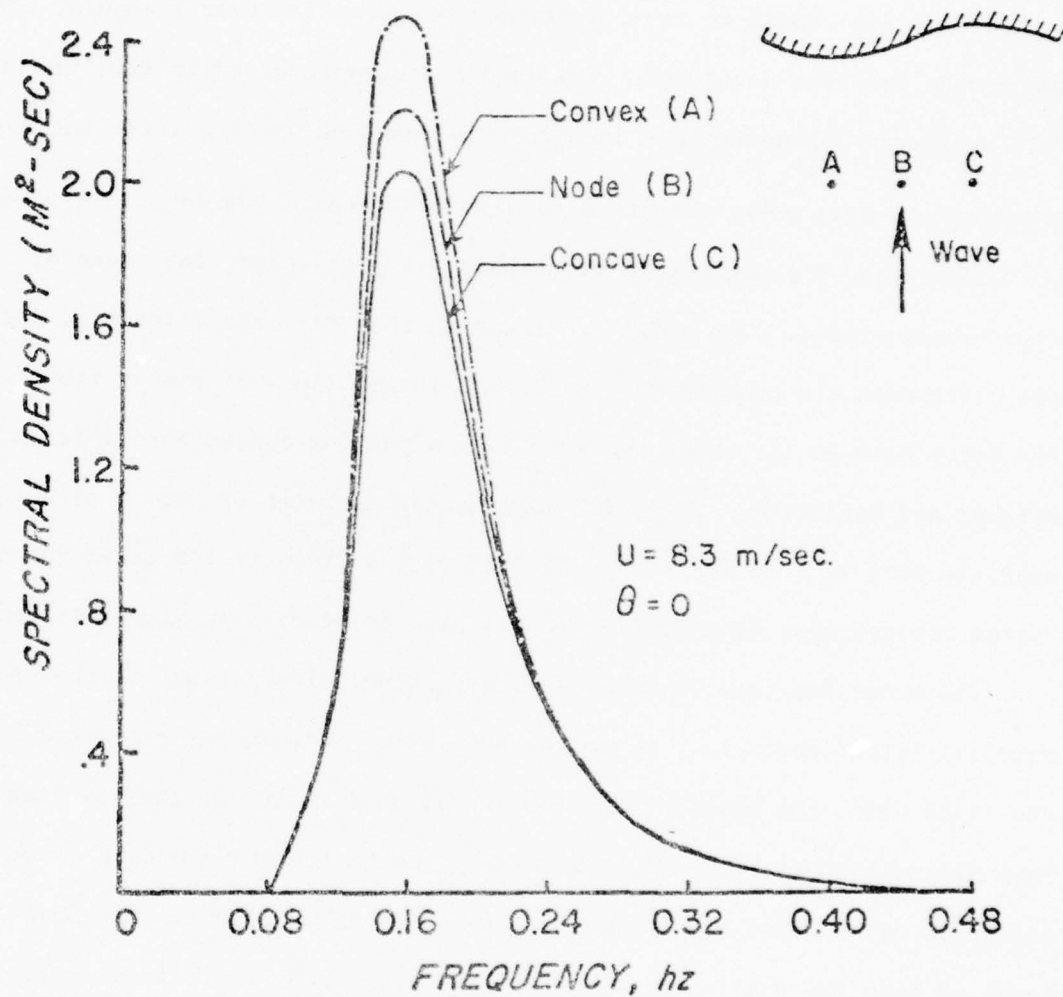


Figure 9 Illustrations of Spatial Variations of Wave Spectrum Over Rhythmic Topography

CHAPTER VI CONCLUSIONS

A numerical model on wave spectrum transformation over irregular bottom topography has been developed. This model proves to be quite effective in predicting wave spectra at predetermined locations in shoaling water outside the surf zone provided the deep water wave condition is specified. The numerical scheme has the advantages of being simple and stable, in particular, for cases of relatively mild bottom slopes. The model is linear in that only the effects of shoaling and refraction are considered. Nonlinear interactions of energy transfer from one frequency band to the other and energy dissipations due to bottom friction and turbulence are neglected. Thus, its application to water of very shallow depth or near the surf line is doubtful. The accuracy of results for cases where the bottom topographies vary rapidly at the side boundaries is also untested.

The model has been verified against analytical solution for the case of parallel bottom contours. As to the comparison between the predicted result and the field data, the agreement is considered good except at the low frequency end. This discrepancy is believed to be mainly due to the assumed deep water input wave condition rather than the nonlinear effects mentioned above. Simultaneous measurements of deep and shallow wave conditions, preferably directional spectrum, should be conducted to establish the field applicability of the model.

It should be noted here that since the present numerical scheme deals with wave height rather than spectral energy density as the basic variables in the governing equations, the incorporation of frictional effects and turbulence energy dissipation is a rather simple matter if such dissipations can be empirically related to the diminishing of wave height. Furthermore, the interactions between waves and currents could also be included as demonstrated by Noda.

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APPENDIX

Computer Flow Chart and
Fortran Listing for Spectrum
Transformation

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C      A PROGRAM TO COMPUTE SPECTRUM TRANSFORMATION OVER IRREGULAR
C      BOTTOM TOPOGRAPHY
C      GS IS ARRAY OF GRAIN SIZE
C      DW IS ARRAY OF DEPTH OF WATER, XBAR IS ARRAY OF SEA LEVEL INTERCEP
C      MI = MAX VALUE FOR I SUBSCRIPT, NOT TO EXCEED 100
C      MJ = MAX VALUE FOR J SUBSCRIPT, NOT TO EXCEED 100
C      TITL = IDENTIFYING TITLE FOR EACH SET. (12A6).
C      NORAYS = NO. RAYS IN EACH SET. (13).
C      T = WAVE PERIOD, SECS, (F8.2).
C      H = DEEP WATER WAVE HEIGHT, (F6.2).
C      X,Y = STARTING COORDINATES. (F7.2).
C      A = INITIAL DIRECTION OF RAY. (F7.2).
C
COMMON DW(50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20),C(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N,IT
COMMON/STRESS/SIGXY(50,20)
DIMENSION QSX(100,100),QSY(100,100)
DIMENSION GS(50,20)
DIMENSION TITL(12),FMT(2)
DIMENSION SAVE(50,20)
COMMON/DTI/DTIS(50)
COMMON/SPC/H1(20),SIG(20),E0(20),E1(50,20)
COMMON/DELTW/DDW
INTEGER OPT
97  FORMAT (/// ' THE ORIGINAL MATRIX OF WATER DEPTH IN METER' )
908  FORMAT (///,1X,'TIDAL LEVEL=NN+',F6.2)
98  FORMAT (///,1X,'DX=',F5.2,3X,'DY=',F5.2)
99  FORMAT (////,1X,'NI=',15,3X,'NJ=',15)
100  FORMAT(2I5)
101  FORMAT(10X,7F10.1/(8F10.1))
102  FORMAT(3F5.2)
103  FORMAT(15F 8.3)
C      READ BASIC DATA
A=135.C
NT=5.
C      NT=NO. OF DIVISION OF ONE TIDAL CYCLE, NT MUST BE ODD
C      NNT=NO OF TIDAL CYCLES
C      NTS= NO. OF TIME STEPS
TPER=12.25
NNT=2
DT=TPER/FLOAT(NT-1)
NTS=NI*NNT
G=9.80621
PI=3.1415927
PI2=PI *2.0
RAD=180.C/PI
READ (5,100) NI,NJ
C      NI IS NUMBER OF GRID POINTS PERPENDICULAR TO THE SHORE
C      NJ IS THE NUMBER OF GRID POINTS PARALLEL TO THE SHORE
M=NI

```

```

N=NJ
IF (NI.GT.50.OR.NJ.GT.20) GO TO 13
READ(5,102)DX,DY,A0

DX2=DX *2.0
DY2=DY*2.0
EPS=0.0005
NBOUND=5
JJ=N+2-NBOUND
NP1=N+1
READ(5,110)((DW(I,J),J=NBOUND,JJ),I=1,NI)
110 FORMAT(7F5.1)

WRITE (6,99).NI,NJ
WRITE (6,98) DX,DY
WRITE (6,908) A0
WRITE (6,97)
WRITE (6,110) ((DW(I,J),J=NBOUND,JJ),I=1,NI)
DO 73 I=1,NI
DO 73 J=NBOUND,JJ
73 DW(I,J)=DW(I,J)+A0
DO 151 I=1,NI
DO 151 J=1,NBOUND-1
151 DW(I,J)=DW(I,NBOUND)
DO 150 I=1,NI
DO 150 J=12,NJ+1

150 DW(I,J)=DW(I,JJ)
WRITE (6,108)
WRITE (6,103) ((DW(I,J),J=1,NP1),I=1,M)
108 FORMAT (//,1X,'THE AUGMENTED MATRIX OF WATER DEPTH IN METER')

C NBOUND IS THE NO. OF PARALLEL PROFILES BOUNDING THE STUDY AREA ON EACH
THETA0=A
THETA0=THETA0/RAD
THETA0=PI2-THETA0
TIME=0.

DO 13 K=1,10
SIG(K)=0.1*PI*FLOAT(K)

SIGMA=SIG(K)
T=2.*PI/SIG(K)
SI4=SIG(K)**4
SI5=SI4*SIG(K)
GU=8.3
CONS=0.74*((9.8/GU)**4)
AW=0.7779*EXP(-CONS/SI4)/SI5
EO(K)=0.1945/CONS

DDW=EO(K)/(AW)
HO=2.828*SQRT(EO(K))
DO 141 IT=1,1

CALL REFRAC(THETA0,HO,IT ,NBOUND,JJ,K)
106 CALL LEVEL(NI,NJ,TIME)

141 CONTINUE
13 CONTINUE
STOP
END

```



```

SUBROUTINE REFRAC(THETA0,HH,ITER,NBOUND,JJ,K)
COMMON D (50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20),C(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
COMMON/DTI/DTIS(50)
COMMON/SPC/H1(20),SIG(20),E0(20),E1(50,20)
COMMON/DELW/DDW

M1=M-1
N1=N+1
NP1=N+1
NBM1=NBOUND-1
N2=N+2
CALL DEPTH
10 CALL SNELL(THETA0,HH,NBOUND,JJ)
CALL ANGLE(20)
CALL HEIGHT(10,ITER,K)
RETURN
END

```

```

SUBROUTINE DEPTH
COMMON D (50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),W(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N,IT
MM1=M-1
156 DO 157 I=2,MM1
DO 157 J=2,N+1
A=D(I+1,J)
IF(A.LT.0.0) A=0.0
B=D(I-1,J)
IF(B.LT.0.0) B=0.0
C=D(I,J+1)
IF(C.LT.0.0) C=0.0
E=D(I,J-1)
IF(E.LT.0.0) E=0.0
DDDX(I,J)=(A-B)/(2.0*DX)
157 DDDY(I,J)=(C-E)/(2.0*DY)
DO 100 J=1,N+1
100 DDDX(1,J)=DDDX(2,J)
DO 101 I=1,M
DDDY(I,1)=DDDY(I,2)
101 DDDY(I,N+2)=DDDY(I,N+1)
158 RETURN
END

```



```

SUBROUTINE SNELL(THETA0,HH,NBOUND,JJ)
COMMON D(50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
DO 600 J=1,N+2
30 DO 600 II=1,M
    I=M-II+1
    IF(D(I,J).GT.0.0) GO TO 33
    ANG=PI
    WVHT=0.0
    SS=0.0
    CC=-1.0
    GO TO 43
33 CONTINUE
    CALL WVNUM(D(I,J),0.0,0.0,0.0,0.0,RK,SIGMA)
    AA=RK*D(I,J)
    ANG=ASIN(SIN(THETA0)*TANH(AA))
    ANG=PI-ANG
    ARG=2.0*AA
    SHOAL=SQRT(1.0/(TANH(AA)*(1.0+ARG/SINH(ARG))))
    REF=SQRT(COS(THETA0)/COS(ANG))
    WVHT=HH*SHOAL*REF
    HB=0.12*TANH(AA)*PI2/RK
    IF(WVHT.GT.HB) WVHT=HB
    IN=1
    IF(WVHT.GT.HB) IN=0
    SS=SIN(ANG)
    CC=COS(ANG)
43 CONTINUE
    U(I,J)=0.0
    V(I,J)=0.0
    H(I,J)=WVHT
    IB(I,J)=IN
    Z(I,J)=ANG
    SI(I,J)=SS
    CO(I,J)=CC
600 CONTINUE
700 RETURN
END

```

```

SUBROUTINE GROUP(I,J,DCGDX,DCGDY,FF,RK)
COMMON D (50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20),C3(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
DIMENSION DUDX(50,20),DUDY(50,20),DVDX(50,20),DVDY(50,20),
3 DTDY(50,20),DKDX(50,20),DKDY(50,20),E(50,20),DTDY(50,20)
JJ=J
DEP=D(I,JJ)
COSI=CO(I,J)
SINI=SI(I,J)
DCGDX=C.C
DCGDY=C.C
FF=0.0
CALL WVNUM(DEP,COSI,SINI,U(I,J),V(I,J),RK,A)
TA=TANH(RK*DEP)
HBREAK(I,J)=0.12*PI2*TA/RK
COSH1=COSH(RK*DEP)
SECHSQ=1.0/(COSH1**2)
ARG2=2.0*RK*DEP
SINH2= SINH(ARG2)
COSH2=COSH(ARG2)
SINH2SQ=SINH2**2
E(I,J)=U(I,J)*COSI+V(I,J)*SINI+0.5*A*(1.0+ARG2/SINH2)/RK
EE=E(I,J)
C=SQRT(G*TA/RK)
FF=0.5*(1.0+ARG2/SINH2)
CG(I,J)=FF*C
DUDX(I,J)=(U(I+1,J)-U(I-1,J))/DX2
DUDY(I,J)=(U(I,J+1)-U(I,J-1))/DY2
DVDX(I,J)=(V(I+1,J)-V(I-1,J))/DX2
DVDY(I,J)=(V(I,J+1)-V(I,J-1))/DY2
DTDY(I,J)=(Z(I+1,J)-Z(I-1,J))/DX2
DKDX(I,J)=(Z(I,J+1)-Z(I,J-1))/DY2
DKDX(I,J)=RK*((U(I,J)*SINI-V(I,J)*COSI)*DTDY(I,J)-(COSI*DUDX(I,J)+
1SINI*DVDX(I,J))-A*DDDX(I,J)/SINH2)/EE
DKDY(I,J)=RK*((U(I,J)*SINI-V(I,J)*COSI)*DTDY(I,J)-(COSI*DUDY(I,J)
1+SINI*DVDY(I,J))-A*DDDY(I,J)/SINH2)/EE
P=C*(SINH2-ARG2*COSH2)/SINH2SQ
DKDDX=RK*DDDX(I,JJ)+DEP*DKDX(I,J)
DKDDY=RK*DDDY(I,JJ)+DEP*DKDY(I,J)
Q=0.5*G/(C*RK**2)
C3(I,J)=C
DCDX=Q*(RK*SECHSQ*DKDDX-TA*DKDX(I,J))
DCDY=Q*(RK*SECHSQ*DKDDY-TA*DKDY(I,J))
DCGDX=P*DKDDX+FF*DCDX
DCGDY=P*DKDDY+FF*DCDY
RETURN
END

```

```

SUBROUTINE LEVEL(NI,NJ,TIME)
COMMON DW(50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20),
2,DDDY(50,20),C(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
DIMENSION XBAR(20)
C    LINEAR INTERPOLATION OF SEA LEVEL POSITION

111 FORMAT('NO SEA LEVEL INTERCEPT ALONG LINE', 15)
104  FORMAT(7F6.2)
    REAL INCREX

    DO 2 J=1,NJ
    DO 3 I=1,NI
    IF(DW(I,J).GT.'C.') GO TO 4
3  CONTINUE
    WRITE(6,111)J
    GO TO 2
4  INCREX=-DX*DW(I,J)/(DW(I-1,J)-DW(I,J))
    XBAR(J)=FLOAT(I)*DX-DX/2.-INCREX
2  CONTINUE
    RETURN
    END

SUBROUTINE NEWHT(I,J,IFLAG)
COMMON D (50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20),
2,DDDY(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
COMMON/XPR/XBREAK(20),SAVE(50,20)
SAVE(M,J)=H(M,J)
IB(I,J)=1
N1=N+1
N2=N+2
CC1=(V(I,J)+CG(I,J)*SI(I,J))/DY
CC2=(U(I,J)+CG(I,J)*CO(I,J))/DX
HNEW=(CC1*H(I,J-1)-CC2*H(I+1,J))/(CC1-CC2-S(I,J)/2.0)
IF(HNEW.LT.'C.0') HNEW=C.0
SAVE(I,J)=HNEW
IF(HNEW.LE.HBREAK(I,J)) GO TO 850
3  HNEW=HBREAK(I,J)
    IB(I,J)=C
850 CONTINUE
860 IF(ABS(HNEW-H(I,J)).GT.(EPS*ABS(HNEW))) IFLAG=0
    H(I,J)=HNEW
899 RETURN
    END

```

```

SUBROUTINE HEIGHT(ITMAX, KK, K)
COMMON D (50,20), U(50,20), V(50,20), Z(50,20), SI(50,20), CO(50,20),
1H(50,20), CG(50,20), S(50,20), HBREAK(50,20), IB(50,20), DDDX(50,20)
2, DDDY(50,20), C(50,20)
COMMON/CON/G, PI, PI2, RAD, EPS, DX, DY, DX2, DY2, T, SIGMA, M, N
COMMON/STRESS/SIGXY(50,20)
COMMON/XER/XBREAK(20), SAVE(50,20)
COMMON/CCIJ/CC(50,20)
COMMON/SPC/H1(20), SIG(20), E0(20), E1(50,20)
COMMON/DELW/DDW
DIMENSION AAW(50,20)
DIMENSION AZ(50,20)
COMMON/DTI/DTIS(50)
M1=M-1

N1=N+1
C COMEPUTE VALUES OF S(I,J)
DO 500 I=2,M1
DO 500 J=2,N1
IF(D(I,J).LE.0.0) GO TO 500
CALL GROUP(I,J,DCGDY,DCGDY,FF,RK)
DUDX=(U(I+1,J)-U(I-1,J))/DX2
DUDY=(U(I,J+1)-U(I,J-1))/DY2
DVDX=(V(I+1,J)-V(I-1,J))/DX2
DV DY=(V(I,J+1)-V(I,J-1))/DY2
DTDX=(Z(I+1,J)-Z(I-1,J))/DX2
DTDY=(Z(I,J+1)-Z(I,J-1))/DY2
SS2=SI(I,J)**2
CC2=CO(I,J)**2
SIGXX=(2.0*FF-0.5)*CC2+(FF-0.5)*SS2
SIGYY=(2.0*FF-0.5)*SS2+(FF-0.5)*CC2
TAUXY=FF*SI(I,J)*CO(I,J)
SIGXY(I,J)=TAUXY
S(I,J)=CG(I,J)*(SI(I,J)*DTDX-CO(I,J)*DTDY)-(DUDX+DV DY)-(CO(I,J)*DC
1GDY+SI(I,J)*DCGDY)-(SIGXX*DUDX+TAUXY*DUDY+TAUXY*DV DX+SIGYY*DV DY)
500 CONTINUE
N2=N+2
M2=M-2
DO 510 II=1,M2
I=M-II
DO 580 IT=1,ITMAX
IFLAG=1
DO 520 J=2,N1
IF(D(I,J).LE.0.0) GO TO 520
CALL NEWHT(I,J,IFLAG)
520 CONTINUE
IF(IFLAG.EQ.1) GO TO 570
580 CONTINUE
WRITE(6,540) I,IT
540 FORMAT(' RELAXATION FOR THE WAVE HEIGHT FAILED TO CONVERGE',
1'ON ROW',I5,'AFTER',I6,'ITERATIONS')
WRITE(6,541)(H(I,J),J=1,N2)
541 FORMAT(' LAST VALUES OF H ARE',/,(10F13.5))
RETURN
570 WRITE(6,542) I,IT
542 FORMAT(10X,' RELAXATION FOR WAVE HEIGHT CONVERGED ON ROW',I6,3X,'A
1FTER',I6,'ITERATIONS')
510 CONTINUE

```

```

        WRITE (6,6052)
6052  FORMAT (//,1X,'WAVE HEIGHT MATRIX'//)
        WRITE(6,652)((H(I,J),J=1,N1),I=1,M)
        WRITE(6,6053)
6053  FORMAT (//,1X,'BREAKING POSITION ("0" MEANS WAVE BREAKING)')
        WRITE(6,653)((IB(I,J),J=1,N1),I=1,M)
        653  FORMAT(15I8)
        652  FORMAT(15F8.3)
        DO 575 J=1,N1
        A=0.
        DO 560 II=1,M2

        I=M-II
        IF(A.GT.0.) GO TO 575
        DXX=0.
        DDXX=HBREAK(I+1,J)-HBREAK(I,J)
        IF(SAVE(I,J)-HBREAK(I,J))560,560,554
554  DDX=((SAVE(I,J)-HBREAK(I,J))/(SAVE(I,J)-SAVE(I+1,J)+DDXX))*DX

557  XBREAK(J)=FLOAT(I)*DX-DX/2.+DDX+DXX
        A=XBREAK(J)
        GO TO 560
555  DDX=0.
        XBREAK(J)=FLOAT(I)*DX-DX/2.+DDX+DXX
        A=XBREAK(J)
560  CONTINUE
575  CONTINUE
        NBOUND=5
        NB=NBOUND
        JJ=11
        DO 702 I=1,M
        DO 702 J=NB,JJ
        E1(I,J)=H(I,J)**2/8.
702  AAW(I,J)=E1(I,J)/(DDW)
        WRITE (6,112) SIGMA
112  FORMAT (//,1X,'OFFSHORE ENERGY DENSITY AT SIGMA=',F8.3)
        WRITE (6,113) EO(K),AW
113  FORMAT (10X,'E'(K)=',F10.3,10X,'AW=',F10.3)
        WRITE(6,704)SIGMA
704  FORMAT(//,1X,'NEARSHORE ENERGY DENSITY AT SIGMA=',F8.3)
        WRITE(6,651)((AAW(I,J),J=NB,JJ),I=1,M)
651  FORMAT(7F10.4)

        WRITE (6,6051) SIGMA
6051  FORMAT (//,1X,'END OF CALCULATION AT SIGMA=',F8.3,/,
1* ***      ***      ***      ***      ***      ***      ***      ***      ***      ***')
        RETURN
        END

```



```

SUBROUTINE ANGLE(ITMAX)
COMMON D (50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
N1=N+1
N2=N+2
M1=M-1
M2=M-2
DO 200 IT=1,ITMAX
IFLAG=1
DO 210 J=2,N1
DO 211 II=2,M1
I=M-II+1
IF(D(I,J).LE.0.0.OR.D(I-1,J).LE.0.0) GO TO 211
CALL NEWANG(I,J,IFLAG,2)
211 CONTINUE
210 CONTINUE
IF(IFLAG.EQ.1) GO TO 250
200 CONTINUE
WRITE(6,220) ITMAX
220 FORMAT(' RELAXATION FOR THETA FAILED AFTER',I6,3X,'ITERATIONS')
250 WRITE(6,251) I,IT
251 FORMAT(10X,' RELAXATION FOR      THETA      CONVERGED ON ROW',I6,3X,'A
1FTER',I6,' ITERATIONS')
WRITE (6,2501)
2501 FORMAT (//,1X,'WAVE APPROACHING ANGLE IN RADIAN')
WRITE(6,252)((Z(I,J),J=1,N1),I=1,M)
252 FORMAT(15F8.3)
RETURN
END

```



```

SUBROUTINE NEWANG(I,J,IFLAG,IM)
COMMON D(50,20),U(50,20),V(50,20),Z(50,20),SI(50,20),CO(50,20),
1H(50,20),CG(50,20),S(50,20),HBREAK(50,20),IB(50,20),DDDX(50,20)
2,DDDY(50,20)
COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N
C STATEMENT FUNCTIONS
DUDX(I,J)=(U(I+1,J)-U(I-1,J))/DX2
DUDY(I,J)=(U(I,J+1)-U(I,J-1))/DY2
DVDX(I,J)=(V(I+1,J)-V(I-1,J))/DX2
DVDY(I,J)=(V(I,J+1)-V(I,J-1))/DY2
F(I,J)=U(I,J)*COSI+V(I,J)*SINI + 0.5*A*(1.0 + ARG2/SINH2)/RK
DKDY(I,J)=(-(COSI*DUDY(I,J) + SINI*DVDY(I,J))-A*DDDY(I,J))/SINH2
1/FF
DKDX(I,J)=(-(COSI*DUDX(I,J)+SINI*DVDX(I,J))-A*DDDX(I,J))/SINH2
1F
FAC(I,J)=U(I,J)*SINI-V(I,J)*COSI
COSI=CO(I,J)
SINI=SI(I,J)
JJ=J
CALL WVNUM(D(I,JJ),COSI,SINI,U(I,J),V(I,J),RK,A)
ARG2=2.0*RK*D(I,JJ)
SINH2= SINH(ARG2)
FF=F(I,J)
IF(FF.GT.0.0) GO TO 450
WRITE(6,451) I,J,D(I,JJ),COSI,SINI,U(I,J),V(I,J),RK,A
451 FORMAT(10X,'FF IS NEGATIVE--OUTPUT I,J,D,COSI,SINI,U,V,RK,A',
12I5,7E13.4)
RETURN
450 FACI=FAC(I,J)
DEN1=(SINI-COSI*FACI/FF)/DY
DEN2=(COSI+SINI*FACI/FF)/DX
DEN=DEN1-DEN2
ZNEW=(COSI*DKDY(I,J)-SINI*DKDX(I,J)+Z(I,J-1)*DEN1-Z(I+1,J)*DEN2)/
1EN
IF(ABS(ZNEW-Z(I,J)).GT.(EPS*ABS(ZNEW))) IFLAG=0
Z(I,J)=ZNEW
CO(I,J)=COS(Z(I,J))
SI(I,J)=SIN(Z(I,J))
RETURN
END

```

```

      SUBROUTINE WVNUM(D,COSI,SINI,U,V,RK,A)
C SUBROUTINE COMPUTES THE WAVE NUMBER INCLUDING WAVE-CURRENT INTERACTION
      COMMON/CON/G,PI,PI2,RAD,EPS,DX,DY,DX2,DY2,T,SIGMA,M,N

      EPSK=0.0005
      RK=PI2/(T*SQRT(G*D))
      DO 100 I=1,50
      A=SIGMA-U*RK*COSI-V*RK*SINI
      A2=A**2
      ARG=RK*D
      F1=EXP(ARG)
      F2=1.0/F1
      SECH=2.0/(F1+F2)
      SECH2=SECH*SECH
      TT=TANH (ARG)
      FK=G*RK*TT-A2
      FFK=G*(ARG*SECH2+TT)+2.0*(U*COSI+V*SINI)*A
      RKNEW=RK-FK/FFK
      IF (ABS(RKNEW-RK).LE.(ABS(EPSK*RKNEW))) GO TO 110
      RK=RKNEW
100 CONTINUE
      WRITE(6,101) I,RK,T,D,U,V
101 FORMAT(' ITERATION FOR K FAILED TO CONVERGE AFTER',I6,3X,5F10.4)
      RETURN
110 RK=RKNEW
      A=SIGMA-U*RK*COSI-V*RK*SINI
      IF(RK.GT.0.0) GO TO 120
      WRITE(6,130) D,COSI,SINI,U,V,RK,A
130 FORMAT(10X,'RK IS NEGATIVE--OUTPUT D,COSI,SINI,U,V,RK,A',/ ,7E12.4)
120 RETURN
      END

```

NJ = 14

0Y=20.00

$$+N_N =$$

THE ORIGINAL MATRIX OF WATER DEPTH IN METER

0.8	0.8	0.6	0.6	0.6	0.6	0.5	0.4	0.4	0.4
0.3	0.4	0.4	0.4	0.4	0.4	0.3	0.2	0.2	0.1
0.2	0.2	0.2	0.2	0.2	0.2	0.0	0.0	0.0	0.0
0.0	0.4	0.3	0.3	0.3	0.3	0.5	0.6	0.7	0.7
0.6	0.8	1.0	1.0	1.1	1.1	1.1	1.1	1.2	1.2
1.1	1.5	1.5	1.8	1.8	1.8	1.6	1.6	1.7	1.7
1.8	2.0	2.2	2.3	2.3	2.0	2.5	2.5	2.3	2.3
2.5	2.5	2.6	2.6	2.6	2.7	2.8	2.8	3.0	3.0
3.0	3.0	3.2	3.2	3.0	3.5	3.6	3.6	3.5	3.5
3.4	3.5	3.5	3.5	3.5	4.0	4.0	4.0	4.0	4.0
3.5	3.8	3.8	3.8	3.8	4.2	4.1	4.1	4.2	4.2
3.6	3.9	3.9	4.0	4.0	4.5	4.3	4.3	4.3	4.3
3.8	4.0	4.0	4.3	4.3	4.6	4.4	4.4	4.4	4.4

THE AUGMENTED MATRIX OF WATER DEPTH IN METER

[illegible]

NEARSHORE ENERGY DENSITY AT SIGMA= 0.628

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0002	0.0002	0.0003	0.0003	0.0004
0.0004	0.0005	0.0007	0.0008	0.0008	0.0009
0.0008	0.0011	0.0013	0.0011	0.0008	0.0011
0.0011	0.0009	0.0011	0.0010	0.0008	0.0010
0.0010	0.0009	0.0010	0.0010	0.0007	0.0009
0.0009	0.0008	0.0009	0.0009	0.0007	0.0008
0.0008	0.0008	0.0008	0.0008	0.0007	0.0008
0.0008	0.0008	0.0008	0.0008	0.0007	0.0007
0.0008	0.0008	0.0008	0.0007	0.0007	0.0007

NEARSHORE ENERGY DENSITY AT SIGMA= 0.942

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0002	0.0000	0.0000	0.0000	0.0002	0.0005
0.0020	0.0020	0.0020	0.0020	0.0107	0.0107
0.0107	0.0477	0.0361	0.0609	0.0609	0.0918
0.0756	0.1095	0.1490	0.1705	0.0899	0.0970
0.1709	0.1417	0.2713	0.1413	0.0836	0.0960
0.1783	0.1291	0.2142	0.1594	0.0885	0.1137
0.1556	0.1276	0.1762	0.1642	0.0959	0.1249
0.1454	0.1265	0.1524	0.1584	0.1030	0.1304
0.1395	0.1263	0.1390	0.1469	0.1124	0.1303
0.1357	0.1272	0.1320	0.1371	0.1200	0.1287
0.1325	0.1282	0.1290	0.1310	0.1237	0.1270
0.1291	0.1278	0.1278	0.1261	0.1247	0.1256

NEARSHORE ENERGY DENSITY AT SIGMA= 1.257

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0002	0.0000	0.0000	0.0000	0.0002	0.0004
0.0016	0.0016	0.0016	0.0000	0.0086	0.0086
0.0086	0.0380	0.0288	0.0000	0.0483	0.0000
0.0596	0.0855	0.1152	0.0006	0.0986	0.0000
0.1314	0.0735	0.2044	0.0211	0.0427	0.0000
0.1150	0.0760	0.2187	0.0653	0.0395	0.0000
0.1077	0.0809	0.1610	0.1135	0.0424	0.0211
0.1061	0.0856	0.1250	0.1316	0.0539	0.0648
0.1051	0.0908	0.1083	0.1192	0.0729	0.0909
0.1035	0.0958	0.1011	0.1095	0.0882	0.0997
0.1012	0.0987	0.0991	0.1030	0.0964	0.1002
0.0989	0.0989	0.0989	0.0990	0.0992	0.0990

NEARSHORE ENERGY DENSITY AT SIGMA= 1.571

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0001	0.0000	0.0000	0.0000	0.0000	0.0000
0.0007	0.0007	0.0000	0.0000	0.0000	0.0000
0.0039	0.0167	0.0000	0.0000	0.0000	0.0000
0.0259	0.0366	0.0007	0.0000	0.0000	0.0000
0.0267	0.0576	0.0000	0.0000	0.0014	0.0000
0.0297	0.0463	0.1393	0.0092	0.0031	0.0006
0.0364	0.0409	0.1729	0.0044	0.0000	0.0046
0.0428	0.0402	0.0666	0.0000	0.0425	0.0028
0.0468	0.0424	0.0513	0.0905	0.0157	0.0287
0.0482	0.0456	0.0475	0.0573	0.0368	0.0413
0.0478	0.0475	0.0471	0.0508	0.0461	0.0485
0.0472	0.0476	0.0476	0.0483	0.0491	0.0486

NEARSHORE ENERGY DENSITY AT SIGMA= 1.885

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0003	0.0000	0.0000	0.0000	0.0000
0.0000	0.0072	0.0000	0.0000	0.0000	0.0000
0.0000	0.0152	0.0000	0.0000	0.0000	0.0000
0.0000	0.0326	0.0000	0.0000	0.0000	0.0000
0.0005	0.0458	0.0510	0.0000	0.0000	0.0000
0.0054	0.0586	0.0193	0.0000	0.0000	0.0000
0.0128	0.0343	0.0187	0.0015	0.0000	0.0000
0.0193	0.0258	0.0206	0.0803	0.0000	0.0000
0.0226	0.0240	0.0219	0.0382	0.0117	0.0546
0.0234	0.0240	0.0232	0.0261	0.0224	0.0229
0.0235	0.0239	0.0239	0.0244	0.0249	0.0246

NEARSHORE ENERGY DENSITY AT SIGMA= 2.199

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0005	0.0213	0.0000	0.0000	0.0003	0.0000
0.0041	0.0242	0.0011	0.0000	0.0005	0.0000
0.0088	0.0186	0.0059	0.0263	0.0007	0.0000
0.0116	0.0140	0.0105	0.0209	0.0056	0.0026
0.0125	0.0131	0.0123	0.0140	0.0119	0.0113
0.0127	0.0128	0.0128	0.0131	0.0133	0.0132

NEARSHORE ENERGY DENSITY AT SIGMA= 2.513

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0002	0.0000	0.0000	0.0000	0.0000	0.0000
0.0009	0.0004	0.0000	0.0000	0.0000	0.0000
0.0019	0.0080	0.0001	0.0000	0.0000	0.0000
0.0036	0.0086	0.0014	0.0000	0.0000	0.0000
0.0055	0.0090	0.0039	0.0000	0.0000	0.0000
0.0067	0.0078	0.0061	0.0000	0.0000	0.0000
0.0070	0.0073	0.0070	0.0000	0.0000	0.0000
0.0071	0.0072	0.0072	0.0000	0.0000	0.0000

NEARSHORE ENERGY DENSITY AT SIGMA= 2.827

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0002	0.0008	0.0006	0.0000	0.0000	0.0000
0.0011	0.0015	0.0018	0.0000	0.0000	0.0000
0.0019	0.0017	0.0024	0.0000	0.0000	0.0000
0.0027	0.0021	0.0025	0.0000	0.0000	0.0000
0.0031	0.0027	0.0031	0.0000	0.0000	0.0000
0.0032	0.0031	0.0031	0.0000	0.0000	0.0000
0.0032	0.0033	0.0031	0.0000	0.0000	0.0000
0.0033	0.0033	0.0032	0.0000	0.0000	0.0000
0.0033	0.0033	0.0033	0.0000	0.0000	0.0000
0.0033	0.0033	0.0033	0.0000	0.0000	0.0000

NEARSHORE ENERGY DENSITY AT SIGMA= 3.142

0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0000	0.0000	0.0000	0.0000	0.0000
0.0000	0.0004	0.0000	0.0000	0.0000	0.0000
0.0000	0.0004	0.0001	0.0000	0.0000	0.0000
0.0009	0.0005	0.0011	0.0000	0.0000	0.0000
0.0012	0.0008	0.0012	0.0000	0.0000	0.0000
0.0013	0.0010	0.0013	0.0000	0.0000	0.0000
0.0013	0.0012	0.0013	0.0000	0.0000	0.0000
0.0013	0.0013	0.0013	0.0000	0.0000	0.0000
0.0013	0.0013	0.0013	0.0000	0.0000	0.0000
0.0013	0.0013	0.0013	0.0000	0.0000	0.0000
0.0013	0.0013	0.0013	0.0000	0.0000	0.0000

ILLUSTRATIVE SAMPLE
COMPUTATIONS

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The computer program has also been converted to suit the ICL 1906 machine at the Technical University at Braunschweig, West Germany. Field work has been conducted at the Island of Sylt in an attempt to verify the model. A set of sample computations are provided here. Data reduction is in progress and the results will be reported separately at a later date.

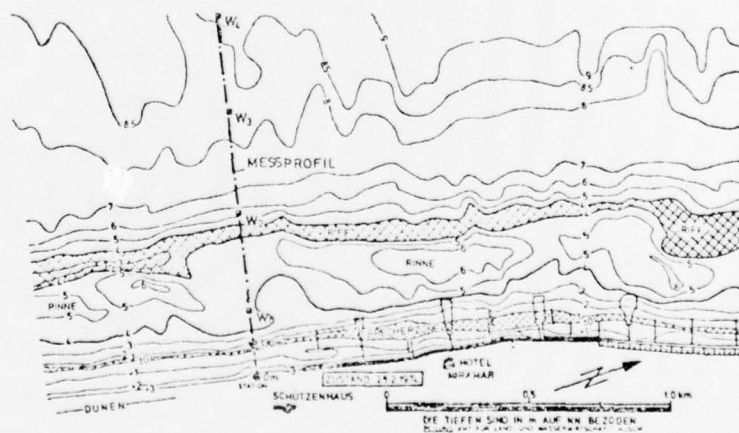
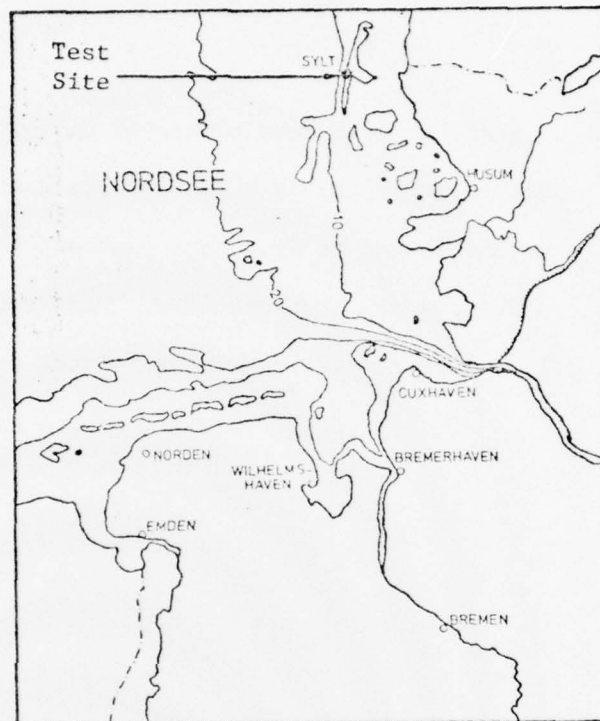
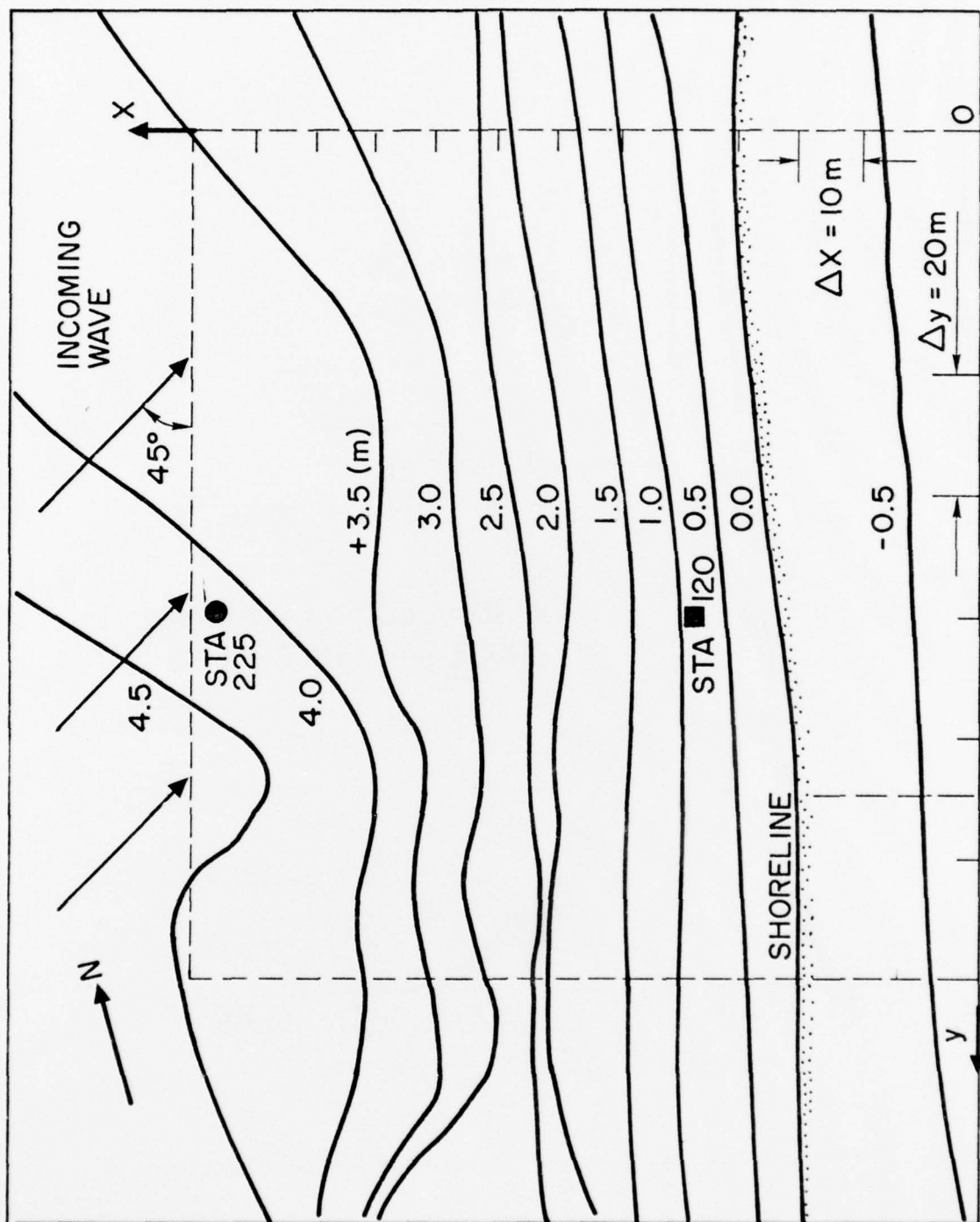


FIGURE 10 Location of Test Site (Island of Sylt/Germany, From Dette, 1974)



MAY 5, 1976

FIGURE 11 Nearshore Hydrograph at Test Site (Island of Sylt/Germany)

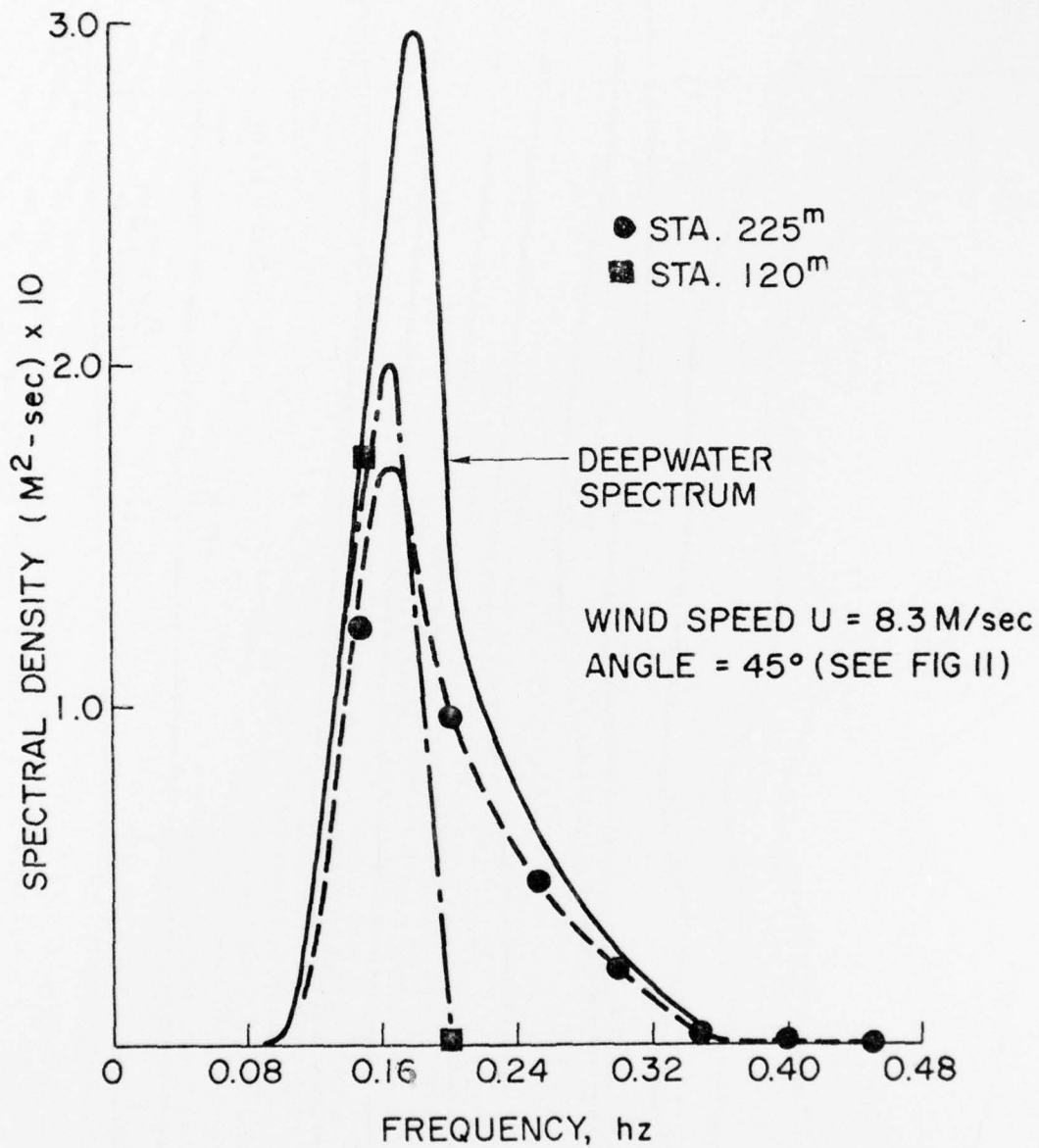


FIGURE 12 Spatial Variations of Wave Spectrum Over Irregular Bottom
 (Island of Sylt/Germany)

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13. ABSTRACT <p>Based on the assumption that the wave energy associated with a narrow frequency band stayed within the band upon refraction, a numerical model based on finite difference method has been developed. This model, utilizing electronic computer, proves to be quite convenient to determine shallow water wave spectra at designated spatial locations of irregular bottom contours provided the deepwater wave spectrum is given. The model has been verified against analytical solution for case of two-dimensional parallel bottom contours and compared in good agreement with field measurement at location well beyond the surf zone.</p>			